## CHAPTER 58

# THE EFFECT OF UNIT WEIGHTS OF ROCK AND FLUID <br> ON THE STABILITY OF RUBBLE MOUND BREAKWATERS 

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## A. ABSTRACT

To study the effect of the specific welghts of armour block material and fluid on the stability of rubble mound breakwaters a total of 110 model tests were made, with varying specific weights of armour and fluid, sizes of blocks and slopes of the breakwater face. The tests indicate that in cases where the specific weights deviate much from usual values, the current design formula (Eq. (1)) should be modified by entering a variabls quantity, $\varphi$, instead of the figure "1" in the denominator. Within the scope of these tests, values varying from $\varphi=0,37$ to $\varphi=1,05$ were indicated. Theoretical considerations seem to show that also hlgher values of $\varphi$ may be expected, in particular on account of the effect of the sloping water surface on the buoyancy of the armour blocks. As neither tests nor analysis havs given conclusive evidence as to under what conditions higher respectively lower values of $\varphi$ should be applied, at present only model tests can give the answer in any particular case.

## B. INTRODUCTION

In locations where suitable rock material is easily available, rubble mound breakwaters with armour blocks of blasted rock, more or less arbitrarily placed, are often economically prsferable. This applies to most breakwaters in Norway, and the stability problems relating to such structures therefore are of particular interest to us.

The specific weights of rock and fluid are important factors in the conditions of stability of such breakwaters. Practically all current design formulae have this form:

$$
\begin{equation*}
Q=\frac{H^{3}}{c^{3}} \frac{\gamma_{r}}{\left(\gamma_{r} / \gamma_{f}-1\right)^{3}}=\frac{H^{3}}{c^{3}} R(\gamma) \tag{1}
\end{equation*}
$$

where $Q$ is the weight of individual blocks neoessary for stability, $H$ is the wave height, $\gamma_{r}$ and $\gamma_{f}$ are the specific welghts of rock and fluid respectively and $c$ is a factor representing all other variables. (In Hudson's well known formula (1), $\mathrm{c}^{3}=\left(\mathrm{K}_{\Delta} \cot \alpha\right.$ ) where $\alpha$ is the slope angle). The above form of the function $R(\gamma)$ has been derived by calculating the buoyancy of the blocks as it would be if the downrushing water were at rest with a horizontal surface. On the whole the experimental evidence in support of the above form of $R(\gamma)$ so far published seems to be somewhat incomplets.

As in most cases the specific welghts do not deviate much from those for which the factor $c$ in Eq. (1) was originally determined, the point is mostly unimportant. But in special cases it may be economically preferable, or even necessary, to use materials with unusual specific weights, in which case the resulting influence on stability becomes a matter of considerable interest.

The purpose of this paper is to present the results of rssearch regarding this problem, carried out during the years 1961-1965 at the River and Harbour* Research Laboratory at the Technical University of Norway, in Trondheim. *) In short, our tests indicated that the figure "1" of Eq. (1) should be replaced by a variable quantity, $\varphi$, which may vary from values of less than one-half to considerably more than one.

The investigation comprises two fairly comprehensive series of model tests, and an attempt at analytical treatment of the process of failure. This attempt has not gaven any full sxplanation of the above indications. In certain respects, however, the analysis has yielded results so far in agreement with experimental data, that its basic concepts may, it would seem, derive some support therefrom.

## C. THE TESTS

Two series of tests were made, Seriss 1, the most comprehensive one, by Mr. Olaf Kydland, in 1961 and 1962, as part of his work for the degree of Licentiatus Technicae, and Series 2 by Mr. Alf T. Sodefjed, in 1965, as part of his work for the degree of Civil Engineer (2), (3) and (4), supervised by the writer.

The scope of the two series of tests is shown in Tables I and II. In Series 1 only one slops of the breakwater model front was ussd, the slope of 1 in 1,5 , most commonly used in actual construction in Norway. Broken natural rock with three different specific weights were used, and in addition blocks broken from a cast of cement and plaster of Paris with sand. For each specific weight parallel tests were made with three different sizes of blocks. An intermediate size was introduced for the 4,52 rock (pyrite), because armour of $105 \mathrm{~cm}^{3}$ blocks of such heavy material could not be broken down in our wave channel.

The blocks were broken manually. Great care was taken to have the blocks of each set of tests as uniform as possible, and at the same time to avold any consistent diffsrencs in shape between blocks of different materials. The weights of individual blocks were kept within $10 \%$ of the average for each group, and the ratio of the greatest to the smallest of the linear dimensions of sach block was kept below 2,5.

Also the specific welghts of the fluid were varied by using, besides fresh water, solutions of NaCl with specific weights of 1,065 and 1,13 .

In Series 2 three different slopes of the breakwater front were used: 1 in $1,25,1$ in 1,5 and 1 in 2. The same types of block material as in Series 1 were used, but in this serles all blocks weighed about the same, which made the volume of any block smaller, the heavier its material. The specific

[^0]welghts varled slightly from those in Series 1 for the four different types. The greatest individual deviation from the average was about $\pm 2 \%$ for the two lighter materials and about $\pm 5 \%$ for the two heavier types. The dimensions of the blocks were kept within the same limits as mentioned for Series 1. Only fresh water was used in this Series.

The model of the breakwater front was built on a wooden slab, Fig. 1 , to eliminate variation in permeability. In Series 1 the slab was covered with two layers of seoondary stones, the mean linear dimension of which varied with the slze of cover blocks from about 1 cm to 3 cm . On top of this sublayer, two layers of the cover blocks, as described above, were placed. In Series 2, a similar arrangement was used, with the sublayer about 5 cm thick.

All tests were made, in an ordinary wave channel, 60 cm wide with depth of water $70 \mathrm{~cm},(14)^{*}$. Each test was started with a wave height well below that causing damage. The height was then raised in decreasing increments as the range of damage was reached.

In Serles 1 the periods were chosen so as to have in all cases as nearly as possible the same steepness of wave at breakdown of the model. In Series 2 the period was $1,8 \mathrm{~s}$ in all tests recorded here.

The wave generator was run continously for 20 mln in Series 1 and 15 mln in Series 2 at each wave height. Secondary reflexion from the padde could not be entirely avoided, but the model was built on wheels as done by Hedar (5), and was moved in each case to a position where the uprush with and without this secondary reflexion were practically equal.

In all tests the degree of damage was noted, as the wave height was increased. The extent of damage was given as the percentage of the total number of cover blocks within a certain specified region, which had rolled down the slope. The wave height corresponding to a given percentage was determined by linear interpolation.

A major problem was how to build all the models sufficiently alike. It proved difficult to avoid a certain improvement of the stability of the model as the routine of the operators improved with time. In fact the first set of tests of Series 1 had to be discarded, because the date on which the model had been built, appeared as a dominant variable in the results.

This difficulty was, it is believed, fairly well overcome in the subsequent tests, in the first place by adopting a strictly standardized method of building the model, and in the second place by making up the test programme so that a possible effeot of improved routine on the average rem sults should be about the same for all combinations tested.

In Series 2 this plan could not be followed throughout, because the tests with slope 1 in 1,5 were decided on only after the other tests were completed. This may partly explain why models with slope 1 in 1,5 apparently were more stable than those with slope 1 in 2 , as seen in Table VI.

The standard method adopted for placing the cover blocks was not quite the same in the two series.

[^1]In Serles 1 the blocks of the first layer were dropped on to the sublayer at a point about twice the expected wave helght for "no damage" above the SWL, and from there rolled down the slope till it stopped against the blocks already placed. If it stopped earlier, the rolling was started again by touching with a finger. The eecond layer of cover blocke was placed in the eame way, but here the finger assistance was more frequently needed, becauee of the greater roughnese of the elope on which the blocks must roll.

In Serles 2 each block was dropped ae directly as possible into its intended place. Blocks belonging below the SWL were dropped from the water surface and those belonging higher up from a helght of some 5-10 cm above the breakwater face. The upper edge of the part of a cover layer already placed was kept sloping from one side of the wave channel to the other, at an angle of about $45^{\circ}$ against the axis of the wave channel when seen normally agalnst the face of the breakwater. Thereby each individual block was guided sideways into a position where it mostly came to rest against two of the blocks previously placed, instead of just one. It was found that this method gave greater stability than that used in Series 1. Thie difference should be more pronounced the steeper the breakwater front is. That may be part of the reason for the relatively low stablity found with a slope of in 2 in this series.

In both Series, at least three identrcal tests were made with each slope, each size of blocks and each combination of specific weights.

In Series 1, If any one of these three tests gave results deviating more than $\pm 10 \%$ from the average of the three, that test was discarded and a new one made. With $3 \times 20=60$ programmed tests, only two individual rem sulte were discarded due to this $10 \%$-rule, whlle three more were diecarded due to other irregularities discovered during the tests, although their results were within the $10 \%$ lamit.

In Series 2, a somewhat stricter rule was used, requiring that the total difference between the maximum and minımum results of identical teets should not exceed $10 \%$ of the average. Here $3 \times 12=36$ tests were programmed, but several tests were repeated more than twice, so that in total 55 tests were made. Of theee 45 gave results within the adopted limit, while 10 fell outside, mostly for obvious reasons.

The wave data pertaining to the tests of the various combinations of specific welghts, sizes of blocks and slopee of the breakwater face summarized in Table I and II, may be seen from Tables III and IV.

## D. PRESENTATION OF RESULTS

The results of the tests are most easlly preeented by bringing Eq. (1) on linear form, and introducing a possibly variable $\varphi$ instead of the fixed quantity, 1. Eq.(1) may be written:

$$
\begin{equation*}
H / k=\lambda=D\left(\gamma_{r} / \gamma_{f}-\varphi\right) \tag{2}
\end{equation*}
$$

Here $Q$ is replaced by $\gamma_{r} \cdot \psi=\gamma_{r} \cdot C_{Y} \cdot k^{3}$, where $\forall 1 s$ the volume and $k$ is a characteristic linear dimension of the cover block in question, $C_{F}$ is a "coefficient of volume" and $D=C_{F} 1 / 3 . c$.

If the general form of the widely accepted Eq. (1) is reasonably correct, aeide from the value of $\varphi$, then the observed values of $\lambda=H / k$ should, when plotted againgt $\gamma_{r} / \gamma_{f}$ as abscissa, group themselves about a straight line, which line will define the values of $\varphi$ and $D$.

In Fig. 3 the average values of $H / k$ fron each set of three parallel tests for each one of the 20 combinations indicated in Table I for Series 1 have been plotted, as observed with $1 \%$ of damage. It is seen that the data are all reasonably close to the straight line drawn in full, which corresponds to $\varphi=0,44$ and $D=0,99$. The data for the heaviest rock material tested, pyrite, fall somewhat below the line, which is more or lese evident throughout both serles of teets. Naturally the drawing of the best fitting straight line may be disputed, but $1 t$ would hardly seem reasonable to draw the line eo ae to bring the value $\varphi$ eloser to 1 than indicated in Fig. 3 .

Actually the line has been drawn after study of simılar diagrams for each of the five groups of combinations tested in Series 1, shown in Figures 4 to 8. In theee diagrams have been plotted the maximum and minamum and the mean value of $\lambda$ found in each of the three individual tests made for each combination, at $1 \%$ of damage. In the same figures the stralght lines corresponding to higher percentages of damage have been shown. For the sake of clarity the data themselves have not been inoluded, but the agreement with the etraight lines is as good as for $1 \%$, or better.

For each of the five groups of combinations, values of $\varphi$ and $D$ corresponding to $1 \%$ and to $4 \%$ of damage have been taken off the diagrame and tabulated in Table V.

The results of the tests of Series 2 for $1 \%$ of damage have been similarly plotted in Fig. 8. As practically the same block welght was used throughout thie eeries, there is just one group of combinations for each value of the angle of slope, $\alpha$. Corresponding values of $\varphi$ and $D$ have been taken off these diagrams and entered in Table VI. Similar diagrams for $10 \%$ of damage have been plotted (not shown) and values of $\varphi$ and $D$ shown in Table VI.

It ie seen from the diagrame and tables that higher values of $\varphi$ are coneistently found for higher percentages of damage, that is for higher stability of the remaining blocks on the breakwater front. Similarly Series 2 gave higher values of $\varphi$ than Series 1 , ae well as higher etability.

During the tests notes were carefully taken of the locations on the elope from which blocks were succeseively washed away. In Fig. 10 is shown how the damage was distributed over the slope, relatively to the wave heights, for all tests of both series.

## E. CONDITION OF STABILITY

A full theoretical explanation of the variation of $\varphi$ indicated by the tests would be most desirable, but the problem is very complicated, and no full eolution, however approximate, has as yet been found. Nevertheless, a rational etady of the oonditions of stability of the armour blocke on a rubble mound breakwater slope, based on fairly reasonable assumptions, may be of some value in clarifying part of the problem.

Any such study must be based on a certain concept of the mode of fallurs of an irregular block of stons forming part of the cover layer on a rubble mound breakwater, as $1 t$ is being washed away by the downrushing water. *) Ths main question is: What will, in most casss, bs the initial morement of such a block ?

Some investigators, among them Svee (7), have assumed that at certain moments some block may become sntirely fres of restraint from neighm bouring blocks and be thrown right out into ths downrushing stream. I have no doubt that this may, and occasionally does occur. It was expressly noted by Kydland, who psrformed with acute observation the tssts of Series 1, that very of ten there sesmed to bs a lockering of the corsr layer around the SWL, bsfore real damage started. Probably some few blocks may thsn have become entirely fres of restraint.

Nevertheless, the writer is inclined to believe that the mode of failurs assumed by Hedar (5), whereby the moving block rolls away, initially in contact with its downstrsam neighbour, corresponds mors nearly to what usually happens. It is hard to ses how a block, once it starts to lift from lts base, can avoid being presssd by the downrushing stream against its neighbour below.

On this basis, and rsferring to the foros diagram in Fig.11, we shall study ths condition of stability of a block "n" against rotation about its point of support, $A_{n}$, on block "n +1" below. Block "p" may, or may not, be stsadied by contact with the block "n-1" above. **) In this Section we shall assume that it is not, ses Section $H$.

If block "n" ls free of contact with block "n-1", its stability depends largsly on the angle $\theta$. Blocks who happen to have the smallsst angle $\theta$ will, other conditions being equal, roll away first.

The forces to be considered are ths weight of the block, Q, its buoyancy, $B$, (which is not directsd vertically and is not equal to $\gamma_{f} \ddagger$ ) a drag force, $F_{D P}$ and an inertial forcs, $F_{M P}$ both expected to act parallel to the slope at soms distance $E k / 2$ above the csnter of gravity of block " $n$ ", and a lift force due to the parallsl vslocity, Fip. Finally there is introduced an hypothetical normal force, $F_{h}$, directed downwards and proportional to the volume of block $n$, not to its projected area. This hypothstical force will be discusssd latsr.

In the "dstailsd summary" previously printed, also a normal drag force due to a suppossd currsnt dirsctsd out of the brsakwatsr body was included. Subsequsnt study has indicated that within ths region close to ths SWL any normal vslocity may bs quits small and may possibly sven bs directed into, not out of the brsakwater body. The assumption of an outward normal drag force of any consequence has thsrefore been droppsd.

[^2]The forces on block "n", Fig. 11, may be written:

$$
\begin{align*}
& Q=C_{\forall} \gamma_{r} k^{3} \\
& B_{N}=C_{\forall} \gamma_{f} k^{3} \cos \alpha \\
& B_{P}=C_{\forall} \gamma_{f} k^{3} \cos \alpha \tan \beta \\
& F_{D P}=C_{A P} k^{2} C_{D P} \gamma_{f} C_{V P} H  \tag{3}\\
& F_{M P}=C_{\forall} k^{3} \gamma_{f} C_{M P} a_{P} / g \\
& F_{L P}=C_{A N} \cdot k^{2} C_{L P} \gamma_{F} C_{V P} H \\
& F_{h}=C_{h} \gamma_{f} C_{\forall} k^{3}
\end{align*}
$$

As stated before, $k$ is a characteristic linear dimension of block " $n$ ". $H$ is the height of the regular waves in the wave channel, and ap is the acceleration of the downrushing stream at block "n". The various coefficients, $C$, will be discussed in Chapter F.

Block "n" will be stable against rotation about point $A_{n}$, Fig.11, if
$Q k / 2 \sin (\theta-\alpha)+B_{p} k / 2 \cos \theta+F_{h} k / 2 \sin \theta \geqq$

$$
\begin{equation*}
F_{D P} k / 2(\varepsilon+\cos \theta)+F_{D M} k / 2(\varepsilon+\cos \theta)+F_{L P} k / 2 \sin \theta \tag{4}
\end{equation*}
$$

By entering equations (3) in Eq. (4) and arranging the terms we
arrive at the following condition of stability of block "n": *)

$$
\begin{equation*}
\frac{k}{H}=\frac{1}{\lambda}=\frac{A_{1} \mu_{1}+A_{2} \mu_{2}}{\gamma r / \gamma_{f}-\left[\psi+C_{M} \mu_{1} a_{P} / g-\mu_{2} C_{h}\right]} \tag{5}
\end{equation*}
$$

where:

$$
\begin{array}{ll}
A_{1}=\frac{C_{A P} C_{D P} C_{V P}}{C_{t}} & A_{2}=\frac{C_{A N} C_{D N} C_{V N}}{C_{*}} \\
\mu_{1}=\frac{\varepsilon+\cos \theta}{\sin (\theta-\alpha)} & \mu_{2}=\frac{\sin \theta}{\sin (\theta-\alpha)} \\
\psi=\frac{\tan \theta-\tan \beta}{\tan \theta-\tan \alpha} &
\end{array}
$$

[^3]
## F. DISCUSSION OF COEPFICIENTS

1. Shape coefficients: The projected area of block " n " in parallel and normal direction is $A P=C_{A P} k^{2}$ and $A_{N}=C_{A N} k^{2}$, respectively. The volume of the block, as defined earlier, is $C H 3$. The values $C_{A P}=C_{A N}=1,0$ and $C_{\text {羊 }}=0,5$ represent faurly well the actual shape of the blocks.
2. Drag coefficients in parallel flow, $C_{D P}$ : Reynold's Number in our
cases is mostly between $10^{4}$ and $10^{5}$ at critical stages and the corresponding value of $C_{D}$ for a smooth sphere, given in current literature, is about 0,4 to 0,5 . From the ordinary Prandtl friction formula Hedar (5) deducted for the waves on a break-water front a boundary resistance corresponding to
$\frac{1}{16\left(\log _{10} \frac{14,8}{k}\right.} z^{2}$, where $z$ is the depth of water above the armour blocks, normally to the slope. Subsequent tests by Andersson (9) have andicated that with very rough slopes the figure 14,8 should be replaced by about 5 . Using this figure, and assuming values of $z$ as found in an earlier study (8) at the critical stage of downrush, values of $C_{D P}$ of 0,3 to 0,4 are found. The value $C_{D P}=0,35$ is chosen for use.
3. Lift coefficient in parallel flow, $C_{\text {IP }}$ : The greater parallel velocity above than below a block will create a lift force. Lattle is known on which to base an assumption as to the size of thas force. From tests on pipes placed on the bottom, the Hydraulic Research Station at Wallingford (10), 1961, reported laft forces from about $3 / 4$ to about $1 / 2$ of the corresponding drag forces, (pp 2 and 3), while Johansson, ( (11, p. 32), reports laft forces up to twice the drag force. Here is assumed $C_{L P}=C_{D P}=0,35$.
4. Coefficient of parallel velocity, $C_{V p}$ : Velocities up and down a
breakwater slope are generally taken to be related to the wave helght by the equation: $v_{P}=\sqrt{C_{V P} \cdot 2 g H}$ Here our concern is with the maxamum velocities around SWL during downrush. From the mathematical model previously presented, (8), it is found that $C_{V P}=0,35$ seems to be a reasonable assumption。
5. Coefficient of mass, $\mathrm{C}_{\mathrm{MP}}$ : For reasons stated in (6), $\mathrm{C}_{\mathrm{MP}}=1,5$ has
been assumed.
Entering values of coefficients 1 to 4 in Eq . (6) we obtain

$$
\begin{aligned}
& A_{1}=\frac{1,00,350,35}{0,5}=0,245 \\
& A_{2}=\frac{1,00,350,35}{0,5}=0,245
\end{aligned}
$$

In the stability condition, $E q \cdot(5), A_{1}$ and $A_{2}$ represent the general conditione, as determined by the general shape of blocke and by the hydraulic relations involved. On the other hand, the factors $\mu_{1}$ and $\mu_{2}$ represent the geometrical stability conditions of those individual blocks which, at the moment considered, are just about to be carried off.

The values assumed for the coefficients can all be disputed, but it is believed that none of them should be considered directly unreasonable. If fair agreement with test result can be shown by applying the same values of $A_{1}$ and $A_{2}$ to all combinations of specific weights, block slzes and slope angles, that might be taken to indicate that our condition of stability may not be too unrealistic.

The geometric stability factore, $\mu_{1}$ and $\mu_{2}$ depend, when the angle of slope, $\alpha$ is given, only on the fraction, $\varepsilon$, and the angle, $\theta$. The former is, of course, unknown, but does not play an important part in the calculations. A value, $E=0,15$ has been used here. Using other values, like 0,10 or 0,20 does not ohange the following argument, it just leads to slightly different "best fit values" of $\theta$.

The same percentage of damage should represent the same stability condition and therefore the same value of $\theta$, lrrespective of specific welghts, sizes of blocks or angels of slope, as long as we are dealing with armour layers that have been constructed alike.

## G. CALCULATION OF $\varphi$

Eq. (5) gives $\lambda$ as a linear function of $\gamma_{r} / \gamma_{f}$, like Eq. (2), provided the other members in the equation are independent of $\gamma_{r} / \gamma_{f}$ :

$$
\begin{equation*}
\lambda=\frac{\gamma_{r} / \gamma_{f}-\left[\psi+C_{M} \mu_{1} a_{\rho} / g-\mu_{2} C_{h}\right]}{A_{1} \mu_{1}+A_{2} \mu_{2}} \tag{9}
\end{equation*}
$$

The two equations (2) and (9) then must be identical:

$$
\begin{align*}
& \varphi=\psi+C_{M} \mu_{1} a_{f} / g-\mu_{2} C_{h}  \tag{10}\\
& \lambda=\frac{\gamma_{r} / \gamma_{f}-\varphi}{A_{1} \mu_{1}+A_{2} \mu_{2}}  \tag{11}\\
& \varphi=\gamma_{r} / \gamma_{f}-\lambda\left(A_{1} \mu_{1}+A_{2} \mu_{2}\right)  \tag{12}\\
& D=\frac{1}{A_{1} \mu_{1}+A_{2} \mu_{2}} \tag{73}
\end{align*}
$$

If the angle $\theta$ is known, $\varphi$ can bs calculated from Eq. (11), with the values of coefficients statsd in Section F. By trial calculation it $1 s$ possible to determine those values of $\Theta$ which will agree most closely with the experimsntal values of $\varphi$, taken off the diagrams in Figures 3 to 9. Such "bsst fit values" of $\theta$ have been calculated for sach series (each method of construction) and for two percentages of damage within each series. The following "best fit values" were found:


Ths values of $\varphi$ and $D$, thus calculated from Equations (11) and (12), have been entered in Tables $V$ and $V I$ for comparison with the experimental values. It is sesn that while there are soms differencss, these are mostly quite small, espscially in view of the fact that Eq. (11) gives $\varphi$ as the differencs between two numbers, ths smaller of which is at least twice the difference.

The general rsquirsment stated at the end of Section $F$ thus is fairly Well well satisfied. Also a higher "best fit valus" of $\theta$ is found for the highsr percentages of damage, and higher valuss for Series 2 than for ths lsss stable models of Ssries 1, all of which agrees with what must be expected.

While this agresment csrtainly is no proof of the correctness of the experimental results and of ths condition of stability arrivsd at, it may possibly be taken as an indication that the results may dsserve a certain degree of confidence.

## H. THE ANGLE $\theta$ AS A PARAMETER OF STABILITY

So far we have assumed that our armour block "n", Fig.11, is not steadisd by any contact with its upstream neighbour, "n - 1". If, however, it is so steadied, a certain force, $P_{n-1}$, acting from block " $n-1^{1 "}$ on block " $n$ " must be included in our stability relations.

It seems reasonable to assume that the set of forces acting on block " $n-1$ " at the moment of critical forces on block " $n$ " near the SWL, will not be very different from the sst acting on block "n". If this is so, the force $P_{n-1}$ may be considsred as compossd of a certain fraction, $p$, of the same forces as those already discussed for block "n", including weight and buoyancy. Bassd on this assumption, calculations havs been made, assuming different values of the fraction, $p$. It has been found that entering such a force $P_{n-1}$ does not materially altsr the calculations, the only effsct beang that the "bsst fit values" of ths angle $\theta$ are lowersd somewhat. For instancs, $p=0,2$ leads to 30 to 50 lower values of $\theta$ than $p=0$.

This means that the stabilizing effect of a foroe $P_{n-1}$ is roughly equivalent to a certain increase in that value of $\theta$ which is necsssary for stabilıty. It appears, thsrefore that the angle, $\theta$, may ussfully be oonsidered as a general parametsr of stablity.

## I. THE "EYPOTHETICAL FORCE", $F_{h}$ •

While fair agreement between the experimental values of $\varphi$ and those calculated from Eq. (11) is easily obtainable, the matter with regard to the other equation for $\varphi, E q$. (10), stands quite differently. The first member, $\psi$, must always be greater than 1. The varıation of $\psi$ with $\theta$ and $\alpha$ is shown in Fig. 11, for $\tan \beta=0,40$, and $1 t$ is seen that in particular with the smaller values of $\theta, \psi$ may easily reach values of 1,4 or more. The second member on the right hand side of Eq. 10 must also be positive, and is not negligible. It seems reasonable to use for the acceleration down the slope the values estimated in (8) for the time when the boundary forces at the SWL pass their maximal value ( (8), Table IV, p. 459) . Assuming for ap $\&$ value of about $0,1 \mathrm{~g}$, with $C_{M P}=1,5$, the second member amounts to about 0,20 for the case of $4 \%$ of damage in Series 1 .

Consequently, unless there $1 s$ a third, negative member, due to our "hypothetical force", $F_{h}$, or other causes, only values of $\varphi$ greater than 1 can satisfy Eq. (10).

It may be of some interest to see, if a force like $F_{h}$ should exist, What must be the value of the "hypothetical coefficient", $\mathrm{C}_{\mathrm{h}}$, to make Eq. (10) agree with the experimental values of $\phi$. Therefore, values of $C_{h}$ have been calculated from Eq. (10) for each of the combinations of specific weights, block slzes and slope angles included in the tests, using in each case the experimental value of $\varphi$. In the calculation of $\psi$, tan $\beta$ has been determaned from Eq. (7) of reference (6).

The values of $C_{h}$ thus determined have been entered in Tables $V$ and VI. It is seen they do not vary much. The mean values of $C_{h}$ and the corresponding standard deviations, $\sigma$, are

For Serles 1, with $1 \%$ of damage: $\bar{c}_{h}=0,525, \quad \sigma=5,7 \%$

(In the last figure, the values for $\cot \alpha=2,0$ have been left out)
Consldering the wide variety of conditions included in the tests, the moderate variation in $C_{h}$ seems remarkable, considering that the individual experimental values of $\varphi$ were used in the calculation.

Still, $1 t$ is possible, although hardly very probable, that the agreement found may be accidental, as $1 t$ has not been shown that a force like $F_{h}$ does actually exist. To enter into Eq. (10) $F_{h}$ must be proportional to the volume of the block. It seems reasonable, then, to look for a regular inertial force, due to an accelerated stream into the breakwater body, or a retarded stream out of it. Attempts at showing the exixtance of such accelerations so far have not succeeded.

It may seem difficult to accept the notion of a force like $F_{h}$, in view of the fact that important normal forces directed out of the breakwater have been observed in several investigations, most clearly, perhaps, by Sigurdsson (13).

It should be noted, however, that we are concerned here with the situation slightly below, but quite close to the SWL, where the bulk of the damage took place in our tests (see Fig. 10), while the great upward normal forces have mainly been observed at points further down the slope, close to the trough between downrushing and oncoming wave. While at the SWL or slightly below, the water surface is at its steepest, further down it flattens out and the slope is even reversed. The great effect of surface slope on the pressure distribution in the fluid (see (8), Eq. (3), p. 448) may well be one cause of a force like $F_{h}$. In fact, while numerical evaluation is difficult, there are indications in several of Slgurdssons diagrams of negative (upward) normal forces close to the SWL at certan stages of the wave cyclus.

Finally, in the highly turbulent and most complicated stream of downrushing water around and over the armour blocks there seems to be ample opportunity for development of forces like $F_{h}$, proportional to the volume of the blocks, although the demonstration of such forces, elther by experiment or by theory may be most difficult.

It is concluded, therefore, that the possibility of a force like $F_{h}$ should not be excluded, as far as the case of armour layers of irregular blocks of blasted rock irregularly placed is concerned. In the case of regularly shaped blocks, regularly placed and even bonded, with an all over more smooth breakwater face, the situation may well be quite different.

## J. PRACTICAL CONSEQUENCES

If the indications of the present study should be proved in the main correct, if it has to be accepted that $\varphi$ may assume values as different from 1 as, say 0,5 and 1,2 , not to go to extremes, such values will have to be taken into consideration in the design of rubble mound breakwaters where the use of material with very unusual specific weights are contemplated.

If any of the current design formulae are employed, the correct value of $\varphi$ should be entered, instead of 1. At the same time, of course, the coefficients of the formulae must be changed so as to give correct block welghts at some usual value of $\gamma_{r}$.

In Table VII an example has been shown, based on Hudson's formula (1) with $K_{\Delta}=3,2$ at $\gamma_{r}=2,65$. It is seen that with values of $\gamma_{r}$ close to normal, the difference is not great, but with value like 3,5 or 2,3 the difference should be taken into account, and with still higher or lower values the difference may be decisive.

There remains, however, the big question, what will be the correct value of $\varphi$ in any particular case. While certain indications can be had from the study here presented, a prediction would be hazardous. Therefore, with unusual specific weights, the only safe procedure at present seems to be to base the design on direct model tests with the materials in question, and with all conditions, ancluding those of building the breakwater,as close to reality as possible.

It is to be hoped that further study of the problem will make safe design recommendations possible.

## K. CONCLUSIONS

1. The tests indicate that it may be advisable to replace the term $\left(\gamma_{r} / \gamma_{f}-1\right)$ in current design formulae for rubble mound breakwaters (Eq. (1)) by $\left(\gamma_{r} / \gamma_{f} \varphi\right)$, where $\varphi 1 s$ a variable quantity. Within the scope of the se tests values of $\varphi$ ranging from 0,37 to 1,05 were found.
2. The tests are belleved to be representative, ae great care was taken to eliminate irrelevant variables and the agreement between the various test results seems quite satisfactory.
3. While no full theoretical explanation of the results is given, an analysis of the stability condition of an armour block on a breakwater slope has yielded results in good agreement with the experimental onee.
4. The assumption of a normsl force directed into the breakwater and proportional to the volume of the block leads to quite consistent resulte as regards the magnitude of such a force which would be required for stabilıty under the varıous test condition.
5. The analysis indicated that values of $\varphi$ exceeding those found in these experiments may well occur.
6. Experiments and analysis both indicate that greater values of $\varphi$ are to be expected, the more stable the placing of the armour blocks has been. Also, $\varphi$ increaeed with increase in cot $\alpha$, within the range of $\cot \alpha=1,25$ to 2,0 .
7. The present inveetigation $1 s$ insufficient to permit definite predictions as to what value of $\varphi$ to expect in particular cases. Therefore, where quite unueual specific weighte occur, it is recommended to resort to model tests in each case.

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Fig. 1. The breakwater model.




Fig. 6. Results of Tests, Series 1, Group III. (Table I) Data from Olaf Kydland (2)


Fig. 7. Results of Tests, Series 1, Group IV. (Table I) Data from Olaf Kydland (2)


Fig. 8. Results of Tests, Series 1, Group V. (Table I) Data from Olaf Kydland (2)

Fig. 9. All Results from Series 2 (Table III).


Fig. 10. Cumulative Distribution of Damage along the Breakwater Face, Series 1 and 2.


Fig. 11. Force Diagram for one Armour Block.


Fig. 12. Variation of $\Psi$ with $\theta$ and $\alpha$.

| TABLE I |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Size of Armour Blocks | Combinations of Specific Weights |  |  |  |  | Group of combinations |
|  |  | Types of Armour Blocks |  |  |  |  |
|  | $\gamma_{f}$ | $\begin{gathered} A \\ \gamma_{r}= \\ 1,83 \\ \hline \end{gathered}$ | $\begin{gathered} B \\ \gamma_{r}= \\ 2,66 \end{gathered}$ | $\begin{gathered} c \\ \gamma_{r}= \\ 3,05 \end{gathered}$ | $\begin{gathered} D \\ \gamma_{r}= \\ 4,52 \end{gathered}$ |  |
| $\forall=11-16 \mathrm{~cm}^{3}$ | 1,00 |  |  |  |  | I |
| $\begin{gathered} 2 \\ \forall=32,7 \mathrm{~cm}^{3} \end{gathered}$ | 1,00 |  |  |  |  |  |
| $\forall=52,0-55,7 \mathrm{~cm}^{3}$ | $\begin{aligned} & 1,00 \\ & 1,065 \\ & 1,13 \end{aligned}$ |  |  |  |  | $\begin{gathered} \text { II } \\ \text { IV } \\ \text { V } \end{gathered}$ |
| $\begin{gathered} 4 \\ \forall=104-106,5 \mathrm{~cm}^{3} \end{gathered}$ | $\begin{aligned} & 1,00 \\ & 1,065 \\ & 1,13 \end{aligned}$ |  |  |  |  | ] III |

[^4]| TABLE III |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
| Types of Armour Blocks |  |  |  |  |  |
| Specific Weights and Sizes |  | A | B | C | D |
| $\gamma_{r}$ | $\mathrm{g} / \mathrm{cm}^{3}$ | 1,725 $\pm 2 \%$ | $2,70 \pm 2,2 \%$ | $3,13 \pm 2 \%$ | $4,72 \pm 5 \%$ |
| $Q_{a v}$ | $g$ | 140,1 | 136 | 143 | 140 |
| $\forall a v$ | $\mathrm{cm}^{3}$ | 81 | 50 | 46 | 30 |
| Slope of break water face | 1:1,25 |  |  | - | $\square$ |
|  | 1:1,5 |  |  |  |  |
|  |  |  |  |  |  |
|  | 1:2,0 |  |  |  |  |

In total $3 \times 4=12$ combinations

1) For $\cot \alpha=1,5, \gamma r=1,86 \mathrm{~g} / \mathrm{cm}^{3}$
table III

Wave Data for Tests of Series 1

| Comblnation of spec. welght and size (TableI) | Group of combinations <br> (TableI) | $\gamma_{r} / \gamma_{f}$ | k <br> cm | At 1\% of damage |  |  | At $4 \%$ of damage |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  |  |  | H <br> cm | $\lambda$ | $\left.\tan \beta^{1}\right)$ | H <br> cm | $\lambda$ | $\tan \beta^{1)}$ |
| A4 | III | I, 83 | 5,97 | 8,0 | 1,34 | 0,462 | 10,0 | 1,67 | 0,419 |
| A3 | II | 1,83 | 4,81 | 7,0 | 1,46 | 0,425 | 7,4 | 1,54 | 0,413 |
| A1 | I | 1,83 | 3,16 | 4,1 | 1,30 | 0,439 | 5,0 | 1,58 | 0,395 |
| A3 - S1 | IV | 1,72 | 4,81 | 6,3 | 1,31 | 0,445 | 7,2 | 1,50 | 0,419 |
| A3 - S2 | V | 1,62 | 4,81 | 5,6 | 1,16 | 0,468 | 6,6 | 1,37 | 0,438 |
| A4-S2 | III | 1,62 | 5,97 | 7,2 | 1,20 | 0,481 | 8, 2 | 1,37 | 0,460 |
| A4 - S1 | III | 1,72 | 5,97 | 7,7 | 1,29 | 0,471 | 9,1 | 1,52 | 0,439 |
| B4 | III | 2,66 | 5,96 | 14, 1 | 2,37 | 0,420 | 16,8 | 2,82 | 0,381 |
| B3 | II | 2,66 | 4,76 | 10,0 | 2,10 | 0,416 | 12,7 | 2,67 | 0,360 |
| B1 | I | 2,66 | 2,83 | 6, 1 | 2,16 | 0,452 | 7,2 | 2,54 | 0,419 |
| B3 - S1 | IV | 2,50 | 4,76 | 9,9 | 2,08 | 0,418 | 11,9 | 2,50 | 0,376 |
| B3 - S2 | V | 2,36 | 4,76 | 9,1 | 1,92 | 0,437 | 10,8 | 2,27 | 0,401 |
| C4 | III | 3,05 | 5,92 | 15,8 | 2,67 | 0,396 | 18,8 | 3,17 | 0,351 |
| C3 | II | 3,05 | 4,78 | 12,0 | 2,51 | 0,372 | 15,1 | 3,15 | 0,309 |
| c3-s2 | V | 2,70 | 4,78 | 11,0 | 2,30 | 0,396 | 13, 2 | 2,77 | 0,351 |
| C1 | I | 3,05 | 2,87 | 6,8 | 2,37 | 0,431 | 8,6 | 2,99 | 0,349 |
| c3 - S1 | IV | 2,87 | 4,78 | 11,7 | 2,47 | 0,378 | 14,2 | 2,97 | 0,327 |
| D3 | II | 4,52 | 4,70 | 16,9 | 3,60 | 0,380 | 22,1 | 4,70 | 0,302 |
| D2 |  | 4,52 | 4,03 | 15,4 | 3,82 | 0,400 | 19,7 | 4,88 | 0,337 |
| D1 | I | 4,52 | 2,83 | 10, 7 | 3,78 | 0,402 | 13,0 | 4,60 | 0,397 |

1) Calculated from (6), Eq. 7

TABLE IV

Wave Data for Tests of serıes 2


1) Calculated from (6), Eq. 7 .

TABLE V, SERIES 1
VALUES OF $\varphi, \mathrm{D}$ IND $\mathrm{C}_{\mathrm{h}}$

| $A_{1}=A_{2}=0,245$ |  |  |  |  | $\mathrm{C}_{\mathrm{MP}} \mathrm{a}_{\mathrm{p}} / \mathrm{g}=0,15$ |  |  | $\cot \alpha=1,5$ |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\gamma_{r / \gamma_{f}}$ | At 1\% of damage |  |  |  |  | At $4 \%$ of damage |  |  |  |  |
|  | From test diagrammes |  | Calculatd with$\mathbf{p}=0, \quad \theta=56^{\circ}, \varepsilon=0,15$ |  |  | From test diagrammes |  | Calculated wath$\mathrm{p}=0, \quad \theta=62^{\circ}, \varepsilon=0,15$ |  |  |
|  | $\varphi$ | D | $\varphi$ | D | $\mathrm{C}_{\mathrm{h}}$ | $\varphi$ | D | $\varphi$ | D | $\mathrm{C}_{\mathrm{h}}$ |
| 1,83 | 0,48 | 1,04 | 0,50 |  | 0,480 | 0,64 | 1,40 | 0,54 |  | 0,407 |
| 1,83 | 0,44 | 0,96 | 0,38 |  | 0,519 | 0,65 | 1,32 | 0,64 |  | 0,404 |
| 1,83 | 0,37 | 0,91 | 0,54 | 1,01 | 0,543 | 0,37 | 1,10 | 0,61 | 1,29 | 0,563 |
| 1,72 | 0,39 | 0,99 | 0,42 |  | 0,530 | 0,56 | 1,29 | 0,56 |  | 0,450 |
| 1,62 | 0,50 | 1,04 | 0,47 |  | 0,468 | 0,58 | 1,29 | 0,56 |  | 0,431 |
| 2,66 | 0,48 | 1,04 | 0,31 |  | 0,503 | 0,64 | 1,40 | 0,48 |  | 0,424 |
| 2,66 | 0,44 | 0,96 | 0,58 |  | 0,524 | 0,65 | 1,32 | 0,59 |  | 0,427 |
| 2,66 | 0,37 | 0,91 | 0,54 | 1,01 | 0,536 | 0,37 | 1,10 | 0,70 | 1,29 | 0,552 |
| 2,50 | 0,39 | 0,99 | 0,44 |  | 0,546 | 0,56 | 1,29 | 0,54 |  | 0,469 |
| 2,36 | 0,50 | 1,04 | 0,46 |  | 0,484 | 0,58 | 1,29 | 0,60 |  | 0,447 |
| 3,05 | 0,48 | 1,04 | 0,40 |  | 0,516 | 0,64 | 1,40 | 0,60 |  | 0,436 |
| 3,05 | 0,44 | 0,96 | 0,56 |  | 0,548 | 0,65 | 1,32 | 0,61 |  | 0,450 |
| 2,70 | 0,50 | 1,04 | 0,42 | 1,01 | 0,508 | 0,58 | 1,29 | 0,55 | 1,29 | 0,469 |
| 3,05 | 0,37 | 0,91 | 0,70 |  | 0,548 | 0,37 | 1,10 | 0,74 |  | 0,583 |
| 2,87 | 0,39 | 0,99 | 0,43 |  | 0,568 | 0,56 | 1,29 | 0,57 |  | 0,490 |
| 4,52 | 0,44 | 0,96 | 0,96 | 1,01 | 0,544 | 0,65 | 1,32 | 0,88 | 1,29 | 0,453 |
| 4,52 | 0,37 | 0,91 | 0,78 |  | 0,564 | 0,37 | 1,10 | 0,96 |  | 0,563 |
| $\begin{array}{ll} \overline{\mathrm{C}}_{\mathrm{h}}=0,525 & \overline{\mathrm{C}}_{\mathrm{h}}=0,472 \\ \sigma= \pm 0,0275 & \sigma= \pm 0,0518 \end{array}$ |  |  |  |  |  |  |  |  |  |  |

TABLE VI, SERIES 2
VALUES OF $\varphi$, D AND $C_{h}$

| $\cot \alpha$ | $\begin{gathered} \gamma_{r} \\ \mathrm{~g} / \mathrm{cm}^{3} \end{gathered}$ | At $1 \%$ of damage |  |  | $C_{M P}{ }^{\text {ap/g }} \mathrm{p}=0$,ge |  | At 10\% of damage |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  | From test diagrammes |  | Calculated with $p=0, \theta=66^{\circ}, \varepsilon=0,15$ |  |  | From test <br> dyagrammes |  | Calculated with $p=0, \theta=73^{\circ}, \varepsilon=0,15$ |  |  |
|  |  | $\varphi$ | D | $\varphi$ | D | $\mathrm{C}_{\mathrm{h}}$ | $\varphi$ | D | $\varphi$ | D | $\mathrm{C}_{\mathrm{h}}$ |
| 1,25 | $\begin{aligned} & 1,725 \\ & 2,70 \\ & 3,13 \\ & 4,72 \end{aligned}$ | 0,58 | 1,31 | $\begin{aligned} & 0,56 \\ & 0,50 \\ & 0,50 \\ & \mathbf{1 , 0 8} \end{aligned}$ | 1,28 | $\begin{aligned} & 0,372 \\ & 0,410 \\ & 0,423 \\ & 0,444 \end{aligned}$ | 0,72 | 1,60 | $\begin{aligned} & 0,73 \\ & 0,78 \\ & 0,74 \\ & 1,52 \end{aligned}$ | 1,65 | $\begin{aligned} & 0,280 \\ & 0,315 \\ & 0,328 \\ & 0,342 \end{aligned}$ |
| 1,50 | $\begin{aligned} & 1,86 \\ & 2,70 \\ & 3,13 \\ & 4,72 \end{aligned}$ | 0,68 | 1,53 | $\begin{aligned} & 0,63 \\ & 0,63 \\ & 0,56 \\ & 0,89 \end{aligned}$ | 1,47 | $\begin{aligned} & 0,357 \\ & 0,390 \\ & 0,409 \\ & 0,441 \end{aligned}$ | 0,80 | 1,84 | $\begin{aligned} & 0,79 \\ & 0,86 \\ & 0,73 \\ & 1,37 \end{aligned}$ | 1,85 | $\begin{aligned} & 0,260 \\ & 0,285 \\ & 0,305 \\ & 0,323 \end{aligned}$ |
| 2,00 | $\begin{aligned} & 1,725 \\ & 2,70 \\ & 3,13 \\ & 4,72 \end{aligned}$ | 0,82 | 1,58 | $\begin{aligned} & 0,93 \\ & 1,07 \\ & 1,05 \\ & 1,23 \end{aligned}$ | 1,77 | $\begin{aligned} & 0,276 \\ & 0,316 \\ & 0,338 \\ & 0,386 \end{aligned}$ | 1,05 | 2,16 | $\begin{aligned} & 1,02 \\ & 1,07 \\ & 0,90 \\ & 1,40 \end{aligned}$ | 2,12 | $\begin{aligned} & 0,076 \\ & 0,113 \\ & 0,137 \\ & 0,163 \end{aligned}$ |
| $\begin{array}{ll} \overline{\mathrm{C}}_{\mathbf{h}}=0,380 & \overline{\mathrm{C}}_{\mathbf{h}}=0,306 \\ \sigma= \pm 0,049 & \sigma= \pm 0,0253 \end{array}$ |  |  |  |  |  |  |  |  |  |  |  |

## TABLE VII

Block lexght Required for $H=6,0 \mathrm{~m}$ and $\cot \alpha=1,5$ with $\varphi=0,5$; 1,0 and 1,2 and varying $\gamma_{r}$, Based on Hudson's Formula:
$\because=\frac{\gamma_{r} H^{3}}{k_{\lambda 1} \cot \alpha\left(\gamma_{r} / \gamma_{f}-1\right)}$ and $K_{\Lambda \varphi}=K_{\Delta 1}\left[\frac{\gamma_{r} / \gamma_{f}-1}{\gamma_{r} / \gamma_{f}-\varphi}\right]^{3} ; \quad K_{\Delta 1}=3,2$

| $\gamma_{r}$ | $\begin{gathered} u_{1}^{t} \\ \text { at } \varphi=1,0 \\ \mathrm{~K}_{\Delta 1}=3,2 \end{gathered}$ | $V_{0,5}^{t}$ at $\varphi=0,5$ |  | $\ell_{1,2}^{\mathrm{t}}$ at $\varphi=1,2$ |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  |  | $\begin{aligned} & Q_{0,5}^{\mathrm{t}} \\ & \mathrm{~K}_{\Delta \varphi}=1,41 \end{aligned}$ | $Q_{3,5}^{t} /{ }_{1}{ }_{1}^{t}$ | $\begin{aligned} & \ell_{1,2}^{\mathrm{t}}, \\ & \mathrm{~K}_{\Delta \varphi}=4,79 \end{aligned}$ | $Q_{1,2}^{\mathrm{t}} / \mathrm{Q}_{1}^{\mathrm{t}}$ |
| 2,0 | 105,0 | 67,0 | 0,64 | 143,0 | 1,37 |
| 2,3 | 54,3 | 45,0 | 0,83 | 61,5 | 1,13 |
| 2,65 | 29,7 | 29,7 | 1,00 | 29,7 | 1,00 |
| 3,0 | 19,1 | 21,7 | 1,14 | 17,8 | 0,93 |
| 3,5 | 11,5 | 14,6 | 1,27 | 9,9 | 0,86 |
| 4,0 | 7,3 | 10,4 | 1,40 | 6,1 | 0,82 |


[^0]:    *) Later referred to as the RHRL.

[^1]:    *) See Fig. 7 of References (14). Flgure 2, therefore, is omitted from this paper

[^2]:    *) Wath the slopss of breakwater front here considered, and aside from occasional "shock forces" from uprushing waves, failure is regularly caussd only by ths downrush, as shown by Hedar (5).
    **) Of course, ths real configuration of blocks is not two-dimsnsional, as in Fig.11, and a block may be held by more than ons downstream and one upstream neighbour. This, howevsr, can not materially alter our reasoning.

[^3]:    *) It is interesting to note that in principle the "Initial Motion Condition" given by Kamphuls, 1966 (12) is Identical with Eq. (5) as far as the hydraulac situations treated are alike.

[^4]:    In total 20 combinations

