

# COMPARISON OF DRAG AND INERTIA COEFFICIENTS FOR A CIRCULAR CYLINDER IN RANDOM WAVES DERIVED FROM DIFFERENT METHODS

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## INTRODUCTION

Wave load is the most important aspect in designing of offshore structures, i.e. offshore wind turbines (OWT), oil platforms, etc. This has been discussed extensively in previous studies e.g. Sarpkaya & Isaacson (1981); Sumer et al. (2006); Heideman et al. (1979) and Davies et al. (1990). Generally, estimations of the wave loads are mainly based on the Morison equation (Morison et al., 1950). By applying the Morison approach, the important parameters are the loading coefficients, i.e. the drag and inertia coefficients. There are available methods to estimate these coefficients, i.e.: the least-squares method with fitting on time domain and/or frequency domain; wave-by-wave fitting the method of moments (Isaacson et al., 1991). The accuracy of methods in estimating the force coefficients ( $C_D$  and  $C_M$ ) are discussed in Isaacson (1991). The first approach seems to be the most accurate.

Heideman et al. (1979) and Davies et al. (1990) determined the drag and inertia coefficients in random wave conditions as a function of Keulegan-Carpenter number ( $KC = UT/D$ , where  $U$  is the horizontal velocity,  $T$  is the wave period and  $D$  is the cylinder diameter), see Figure 1. The drag coefficient is considerably scattered at low  $KC$  number.

## METHODOLOGY

In this paper available methods are used to estimate the force coefficients. Those are the max-min method and least-squares method (simplified by fit on wave-by-wave basis). According to the method of max-min values of wave kinematics are determined only at the crest, trough and still water level for each individual wave. The least-squares method in time domain is simplified by Davies (1990) by applying a wave-by-wave basis.

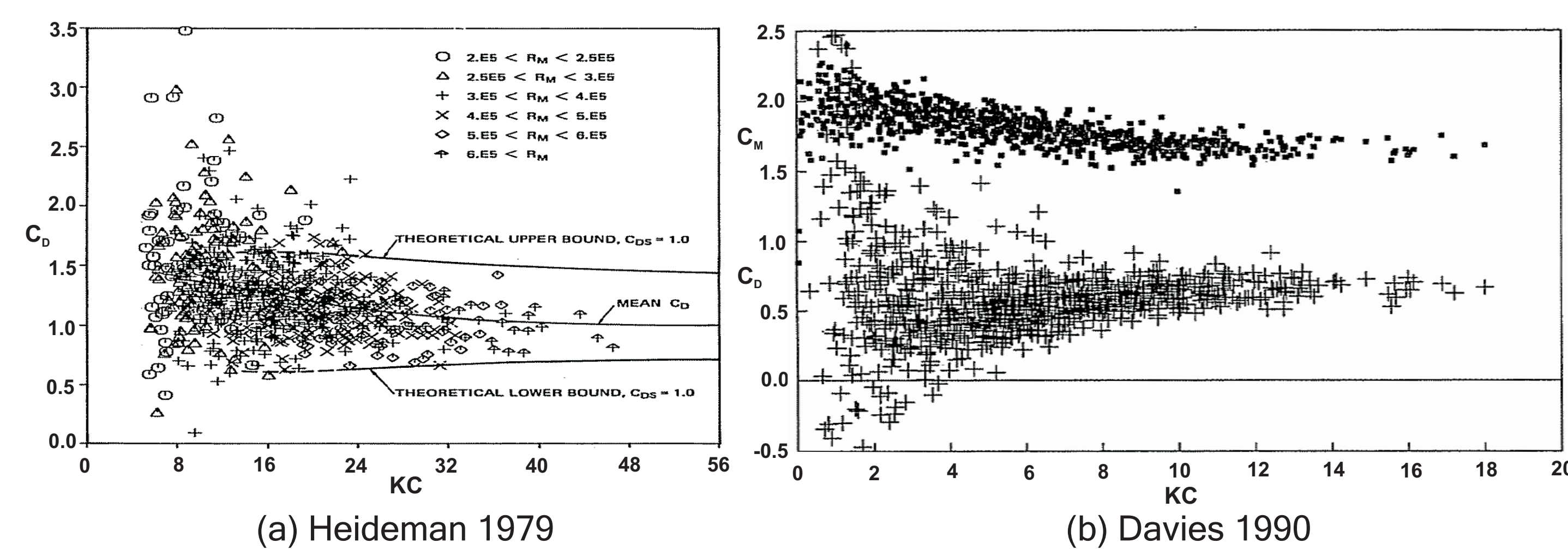


Figure 1: The correlations of the force coefficients and Keulegan-Carpenter number

## STUDY APPROACH

In this study, a combination of analytical approaches and interpreting results of physical model tests is applied. Series of physical model experiments are conducted by fixing a cylinder with a diameter of 30 cm in the wave flume of the Franzius-Institute (Germany), see Figure 2 & 3. Sets of random waves are applied in combination with varying water depths. For each test pressures around the cylinder induced by waves are measured at different elevations. Measured data are used to estimate the coefficients derived from analytical approaches in combination with different wave theories (Airy, Stokes 2<sup>nd</sup> and 5<sup>th</sup> order).

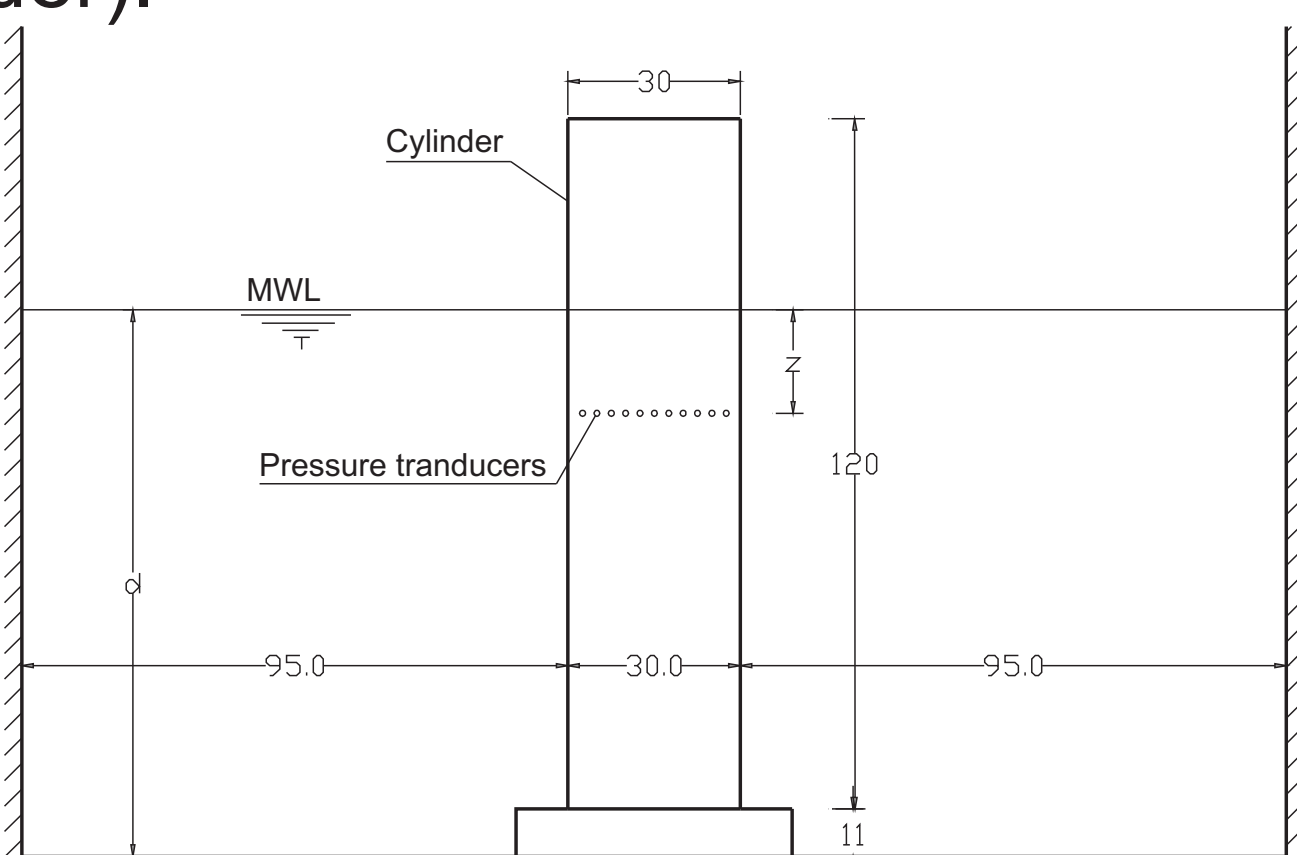


Figure 2: Setting up experiments (units in cm)



Figure 3: Physical model in the wave flume

Morison equation: 
$$F(t) = \frac{1}{2} \cdot \rho \cdot C_D \cdot D \cdot u(t) \cdot |u(t)| + \frac{\pi}{4} \cdot \rho \cdot C_M \cdot D^2 \cdot \ddot{u}(t)$$

Least-squares method: 
$$\begin{cases} f_d \cdot \sum u^4(t) + f_i \cdot \sum u(t) \cdot |u(t)| \cdot \ddot{u}(t) = \sum u(t) \cdot |u(t)| \cdot F_m(t) & f_d = \frac{1}{2} \cdot \rho \cdot C_D \cdot D \\ f_d \cdot \sum u(t) \cdot |u(t)| \cdot \ddot{u}(t) + f_i \cdot \sum \ddot{u}^2(t) = \sum \ddot{u}(t) \cdot F_m(t) & f_i = \frac{\pi}{4} \cdot \rho \cdot C_M \cdot D^2 \end{cases}$$

Max-min method: 
$$C_D = \frac{F_m(t_1)}{\frac{1}{2} \cdot \rho \cdot D \cdot u(t_1) \cdot |u(t_1)|} \quad C_M = \frac{F_m(t_2)}{\frac{\pi}{4} \cdot \rho \cdot D^2 \cdot \ddot{u}(t_2)}$$

## SUMMARY OF THE RESULTS

- Applying different wave theories results in a variation of ~10 % in estimating the wave force coefficients.  
- Relations of drag and/or inertia coefficients, and the Reynolds- and/or  $KC$ -numbers are established. The force coefficients are correlated to the  $KC$  number, however the drag coefficient is rather scattered at low  $KC$  number (Figure 4). This agrees well with previous studies in Davies (1990), Heideman (1979) and Isaacson (1991).

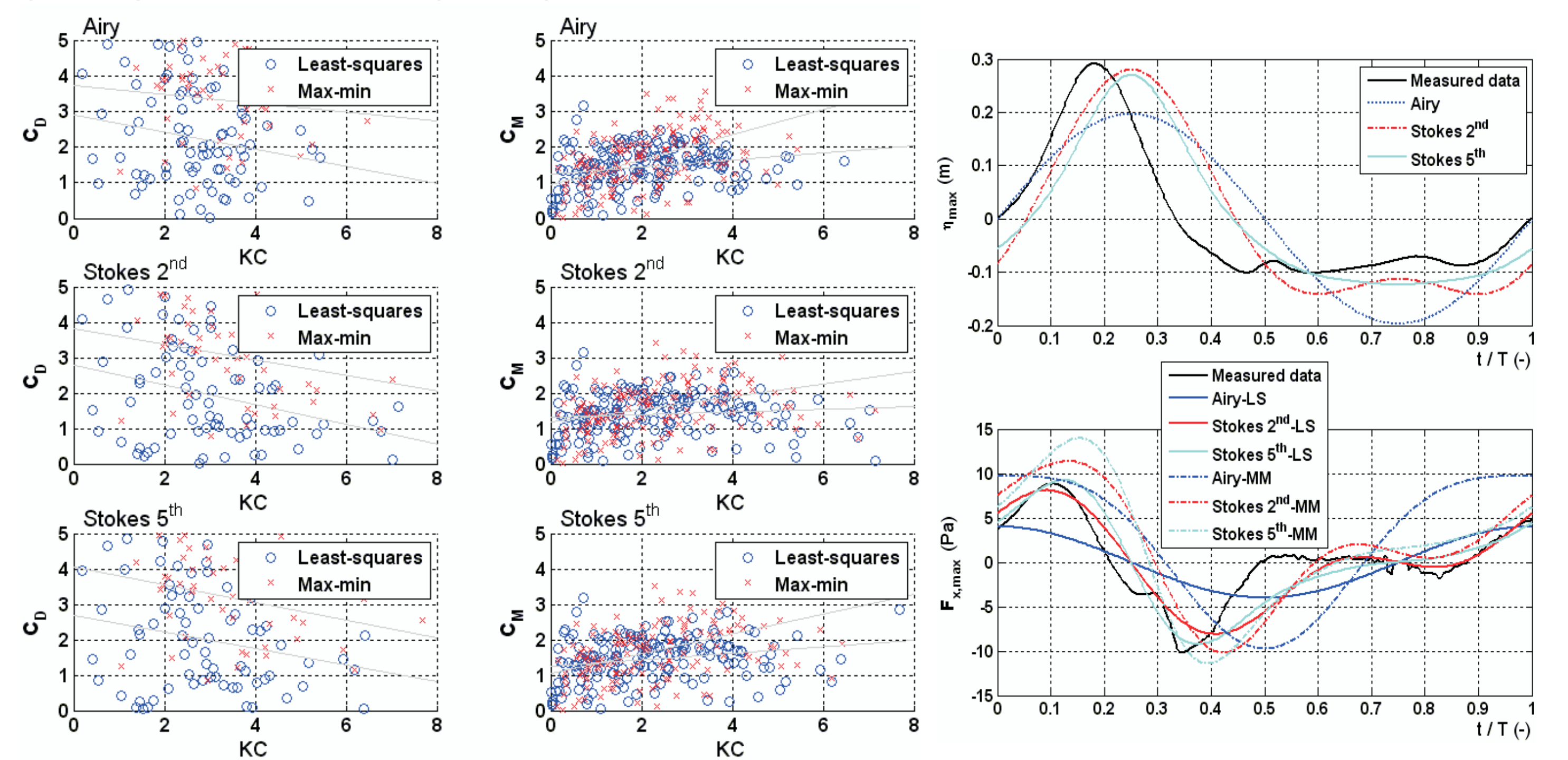


Figure 4:  $C_D$ ,  $C_M$  &  $KC$ -number at level  $z = -0.1$  m

Figure 5: Wave surface & force at  $z = -0.1$  m

Table 1: The mean values of  $C_D$  &  $C_M$  ( $H_s = 0.2$  m,  $T_p = 2.5$  s,  $d = 0.95$  m)

Level	Wave theory	Least-squares		Max-min	
		$C_D$	$C_M$	$C_D$	$C_M$
$z = -0.4$ m	Airy	2.30	1.45	3.21	1.24
	Stokes 2nd	2.03	1.43	3.07	1.21
	Stokes 5th	2.16	1.48	3.16	1.25
$z = -0.1$ m	Airy	2.55	1.38	4.10	1.11
	Stokes 2nd	2.47	1.35	3.85	1.09
	Stokes 5th	2.64	1.42	3.91	1.13

Table 2: The mean values of  $C_D$  &  $C_M$  ( $H_s = 0.25$  m,  $T_p = 3.0$  s,  $d = 0.95$  m)

Level	Wave theory	Least-squares		Max-min	
		$C_D$	$C_M$	$C_D$	$C_M$
$z = -0.4$ m	Airy	2.29	1.51	2.80	1.50
	Stokes 2nd	2.05	1.46	2.52	1.42
	Stokes 5th	2.00	1.53	2.62	1.50
$z = -0.1$ m	Airy	2.56	1.39	3.49	1.24
	Stokes 2nd	2.63	1.34	3.68	1.20
	Stokes 5th	2.66	1.41	3.80	1.27

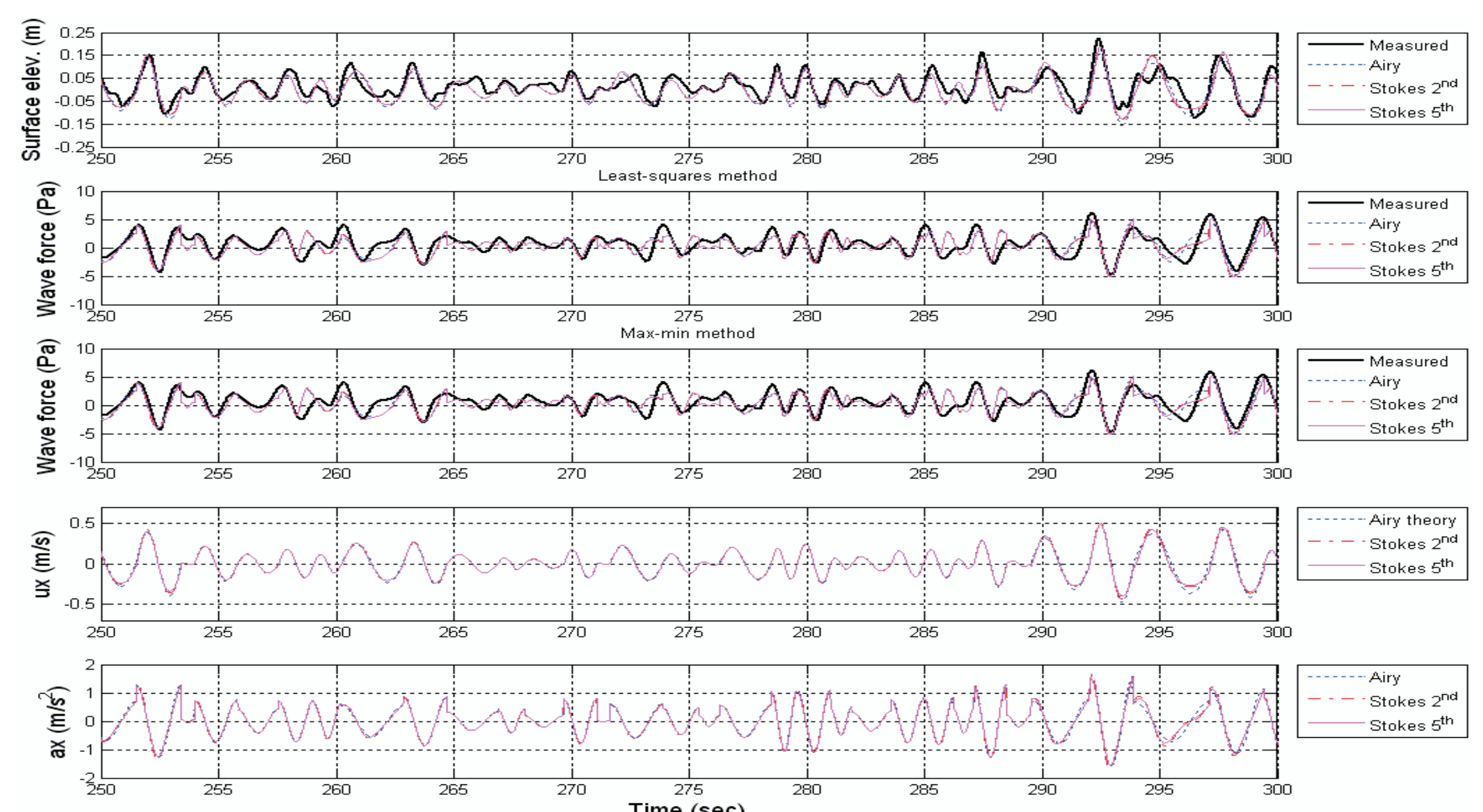


Figure 6: Wave surface, force and kinematics at  $z = -0.4$  m

Figure 5 presents the estimated force at a level of 0.1 m below mean water level (MWL) on the cylinder. It appears that the estimation of Stokes 5<sup>th</sup> order agrees well with the measured data.

Table 1 & 2 show the mean values of  $C_D$  and  $C_M$  at different levels of the cylinder under two wave conditions.

Figure 6 shows the water surface, wave forces and wave kinematics at level of 0.4 m below MWL for a random wave condition.

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