

CHAPTER 122

Sediment Suspension due to Large Scale Eddies in the Surf Zone

by

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ABSTRACT

Laboratory experiments using a fiber-optic LDV system and a small pressure transducer have been made to reveal detailed characteristics of the velocity field in the surf zone and its relationship to the sediment suspension with special reference to the three-dimensional large scale eddies referred to as "obliquely descending eddies", the existence of which was recently revealed by Nadaoka (1986). A conditional sampling technique has been used to find that the obliquely descending eddies bring highly intermittent intensive turbulence to the bottom with the large onshoreward momentum at the upper layer of the water and thus essentially characterize the turbulent flow field in the surf zone. Visual observation and concentration measurements, especially a coherence analysis of two data sets of concentration close to the bottom, have shown that the sediment suspension is mostly governed by such large scale eddies in a wide extent of the surf zone; i.e., the eddies hit the bottom and then lift up the sediment into suspension, yielding the spot-like sediment cloud in accordance with the three-dimensional eddy structure.

1. Introduction

The effect of wave breaking is one of the most significant factors for the coastal sedimentary process. In this respect, Miller (1976), Shibayama and Horikawa (1985), and others have discussed sand suspension by the action of the large scale vortices generated immediately after the wave plunging (Fig. 1). The generation of such vortices, which

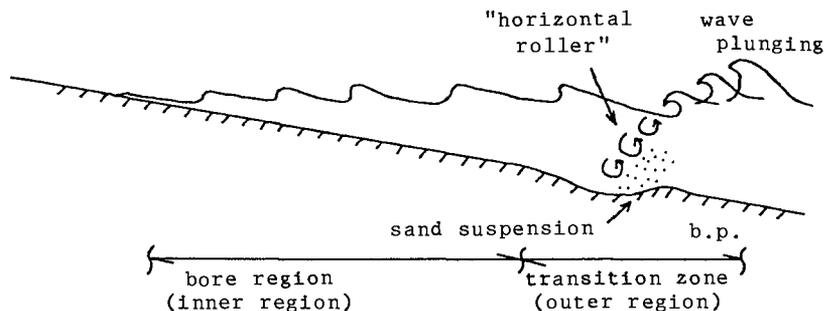


Fig. 1 "Horizontal rollers" and sediment suspension.
(Miller 1976, Shibayama & Horikawa 1986, etc.)

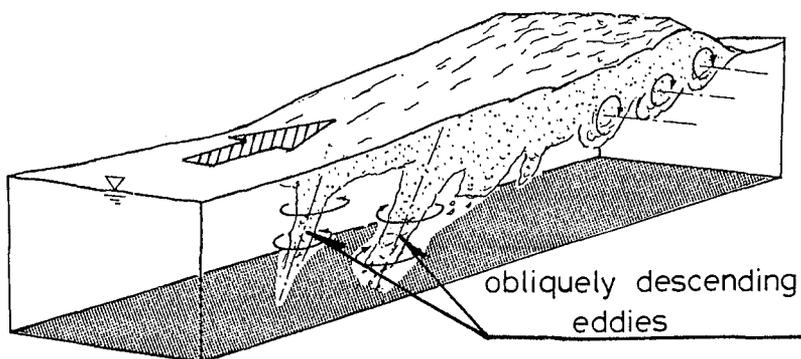


Fig.2 Schematic representation of the large scale eddy structure under breaking waves. (Nadaoka 1986)

are called "horizontal rollers", is a local phenomenon in the sense that it is limited in space within the transition zone or outer zone. In another word, they don't affect sediment suspension in the inner region or bore region.

For spilling breaker condition, on the other hand, Deigaard et al.(1986) have discussed the effect of the turbulence originated from breaking water surface on the vertical diffusion of the bottom sediment, by using the transport equation for turbulent kinetic energy. However, the real physical process of the sediment suspension under spilling breaker as well as turbulent bore has been not yet well understood, mainly because of our limited knowledge of the turbulent velocity field structure within the surf zone.

Recently, Nadaoka(1986) has found the inherent eddy structure of the breaking waves in a wide extent of the surf zone. That is, around the wave crest, dominant eddies have nearly two-dimensional flow structure the axis of which is parallel to the crest line, while behind the wave crest, the

flow structure changes quickly into that of obliquely downward stretched eddies with strong three-dimensionality which is referred to as "obliquely descending eddies" (Fig.2). Therefore, it is worth to pay our attention to these obliquely descending eddies, because they may act as possible agitator of the bottom sediment in the surf zone.

The main purpose of the present study is to clarify the detailed characteristics of the turbulent velocity field in the surf zone and its relationship to the sediment suspension with special reference to the "obliquely descending eddies" under the turbulent bore.

2. Experimental Equipments and Procedure

The wave channel used for the experiments was 0.6 m deep, 0.4 m wide and 20 m long, equipped with a flap-type wave generator. The experiments were made both for a fixed bed and a movable bed. In the former case, a wooden slope of 1 on 20 was installed at one end of the channel and in the latter ground coal of 0.21mm in median diameter and of 1.45 in its specific weight was placed with the depth of 15cm on the slope, as shown in Fig.3. The use of the ground coal for the bed material was to suppress the generation of the ripples under the wave action as much as possible, because in the natural surf zone condition the ripples is hardly developed.

The experiment was carried out under the condition of wave period $T=1.27s$, equivalent offshore wave height $H_0^*=15.5$ cm, offshore wave steepness $H_0^*/L_0=0.061$. Regular waves generated by the wave maker produced a spilling-type breaker, the depth and wave height of which, h_b and H_b , were 16.7 cm and 14.7 cm, respectively. It should be noted that, under the spilling breaker condition, there exists no horizontal roller.

Horizontal and vertical velocity components, u and w , were measured with a two-component fiber-optic laser-Doppler velocimeter of back-scattered fringe mode (FLV, Hino et al., 1984). The water surface fluctuation η was measured with a

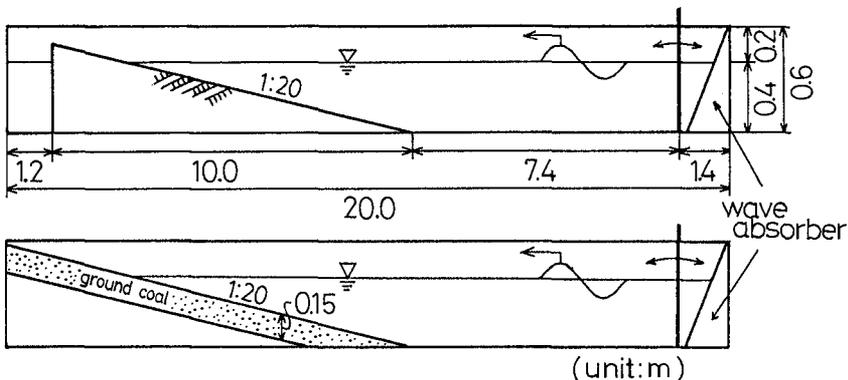


Fig.3 Experimental equipment.

capacitance-type wave gauge at right above the velocity measuring point which was located at $h/h_b = 0.67$ (h is the still water depth). At the measuring section, breaking waves were confirmed to be well developed into so-called "turbulent bore". A small pressure transducer of 1 cm in diameter, 3 mm in thickness and 1 kHz in its frequency response was used for the measurement of bottom pressure fluctuation at the measuring section. The concentration of the suspended bottom sediment for the movable bed experiment was measured by a light-attenuation-type sensor, the position of which was manipulated to be kept at nearly the same height above the bottom through the movable bed experiment.

3. Characteristics of Velocity Field Associated with Oblique Eddies

3.1 Intermittent structure of bottom turbulence and pressure fluctuation

Figure 4 shows the phase-averaged distribution of the turbulence intensity u' and w' at the several heights z above the bottom. The origin of the horizontal axis corresponds to the phase of the zero-up crossing point of the water surface profile. The values indicated were calculated by averaging about 200 records of the high-frequency velocity component data which is obtained through a moving-average filter with a averaging time of 0.1 s. We can see in this figure that the dependence of the turbulence intensity on the phase variation as a whole is not so distinct, though the turbulence intensity at $z = 6.7\text{cm}$ slightly increases around the phase of wave crest. Especially, the turbulence intensity close to the bottom is almost uniform except the region around the phase of the wave crest, where the intensity becomes larger probably corresponding to the increase of the orbital velocity.

Figure 5, on the other hand, is a typical example showing the raw data of the water surface fluctuation η immediately above the velocity measuring point, the bottom pressure fluctuation p , and the velocity fluctuation u and w close to the bottom ($z=0.7\text{cm}$). In this figure, the bottom turbulence occurs not uniformly, but appears in a quite intermittent manner with the bottom pressure disturbances at the instants indicated by the arrows. Further, we should note that such an intermittent occurrence of the bottom turbulence can be found not under the wave crest, but around the phase of the zero-down crossing point. This means that these intermittent turbulence is not caused by the bottom shear stress due to orbital wave motion.

3.2 Conditional sampling analysis of velocity field associated with the oblique eddies

Comparison between Figs. 4 and 5 leads us to the idea that the intermittent characteristics of the bottom turbulence may be lost through the averaging process by the usual phase-averaged method because of the random occurrence of the intensive bottom turbulence, as shown in Fig. 5.

Considering the occurrence of such a highly intermittent

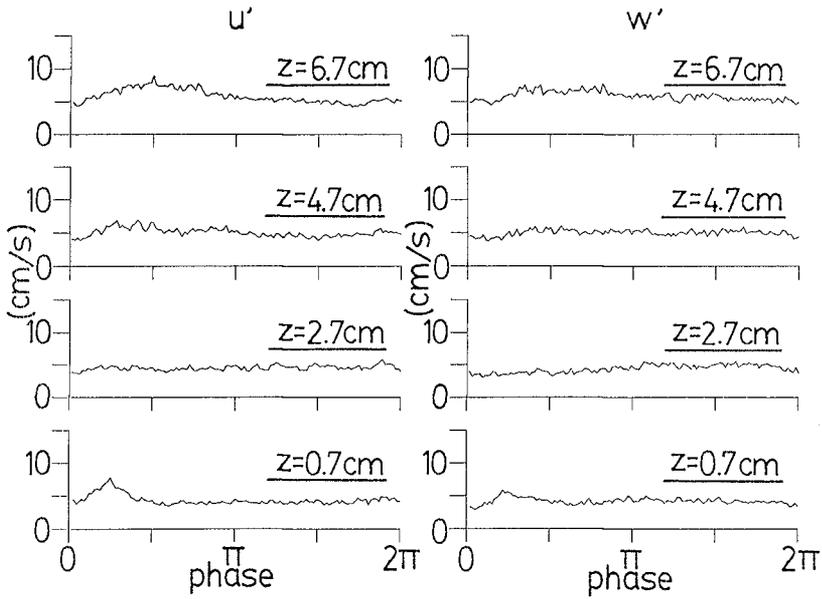


Fig. 4 Distributions of phase-averaged turbulence intensity u' and w' .

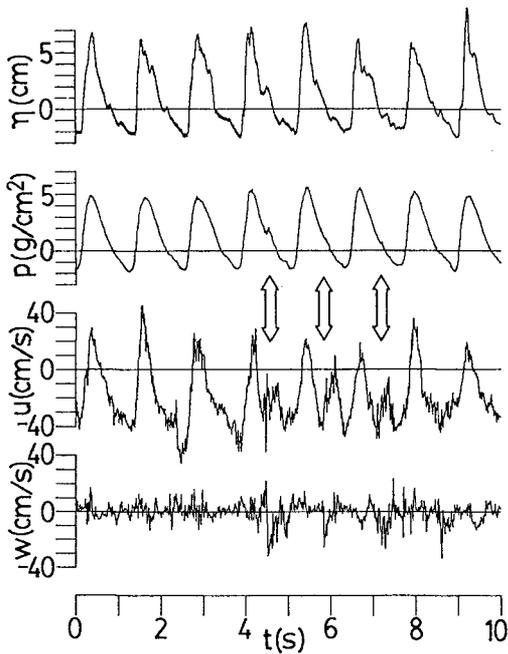


Fig. 5 An example of time series of η , p , u and w .

bottom turbulence behind the wave crest, we can expect that these disturbance is caused by the action of the obliquely descending eddies mentioned first. The results of the visual observation of the eddy evolution seemed to support this expectation. To confirm this point more quantitatively, we have investigated the flow structure related to the intermittent bottom disturbances by applying a conditional sampling technique.

In Fig. 5, we can find also that the pressure disturbance occurs not so frequently as compared with the velocity disturbances. Its possible explanation is that the pressure disturbance occurs only at the instant when the pressure gage is directly hit by an obliquely descending eddy. Considering this characteristics of the pressure record, we can utilize it as a trigger signal for the conditional sampling to detect distinctively the velocity data corresponding to the intermittent bottom turbulence. Namely, by inspecting whether the magnitude of the pressure disturbance exceed the prescribed threshold level or not, we can classify each individual wave record separated by the zero-up crossing method into conditionally sampled or rejected records. The scanning of the pressure disturbance for the conditional sampling was done within the phase interval $[\pi/2, 3\pi/2]$ of each individual wave record, because the apparent pressure fluctuation calculated by the moving-average method may emerge in the phase intervals $[0, \pi/2]$ and $[3\pi/2, 2\pi]$ through the effects of the water surface fluctuation around the wave crest and of the higher harmonics due to the wave nonlinearity (Fig.6). The number of the wave records obtained through this conditional sampling is about 50, 1/10 of the total wave records.

Figure 7 represents the phase-averaged distribution of the turbulence intensity u' and w' , where the solid and broken lines correspond to the conditionally sampled and rejected data, respectively. We can clearly see in this figure that the turbulence intensity of the sampled data at the most upper measuring point attains relatively large value around the wave crest and this intensive turbulence propagates obliquely downward as indicated by the arrows, while the turbulence intensity of the rejected data exhibits no appreciable dependence on the phase variation. This means that the obliquely descending eddies bring intensive turbulence to the bottom behind the wave crest.

Figure 8 shows another result of the conditional sampling analysis where the phase-averaged distributions of the mean velocity, $\langle u \rangle$ and $\langle w \rangle$, of the conditionally sampled and non-conditioned data are represented by the full and broken lines, respectively. In this figure, we can recognize that, in the large turbulence regions, both the onshoreward and the downward mean velocities are enhanced in the sampled records. This fact is directly related with the increase of the downward transport of the onshoreward momentum as described in what follows.

The left and right figures in Fig. 9 represent the phase-averaged distributions of the momentum transport for the sampled and rejected records, respectively. The solid lines indicate the phase-averaged product of the horizontal and the vertical velocity fluctuation, $\langle (\bar{u}-\bar{u}) \cdot (\bar{w}-\bar{w}) \rangle$, where the over bar means an operator to take a time average. Comparing

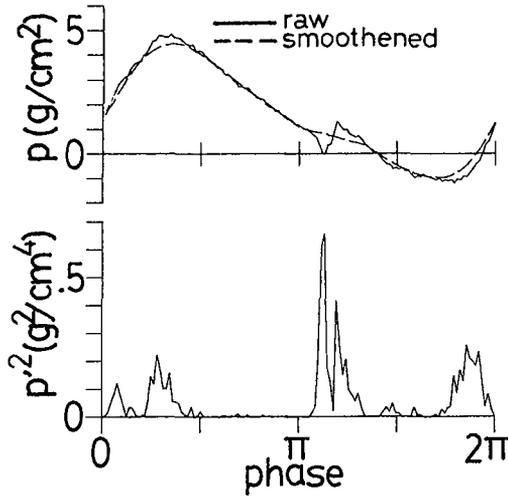


Fig. 6 Raw and smoothed data of p . p' is the difference between these two data.

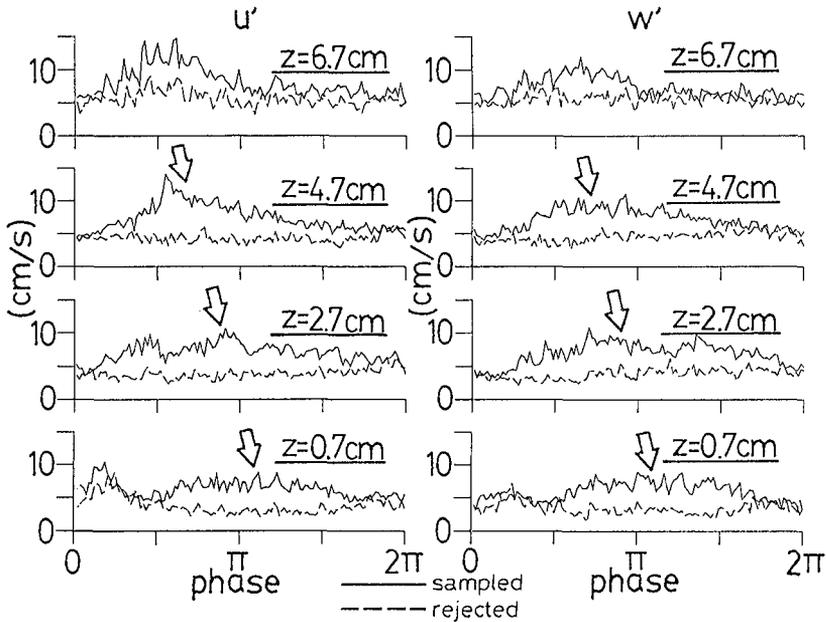


Fig. 7 Phase-averaged distributions of the turbulence intensity u' and w' for conditionally sampled and rejected records.

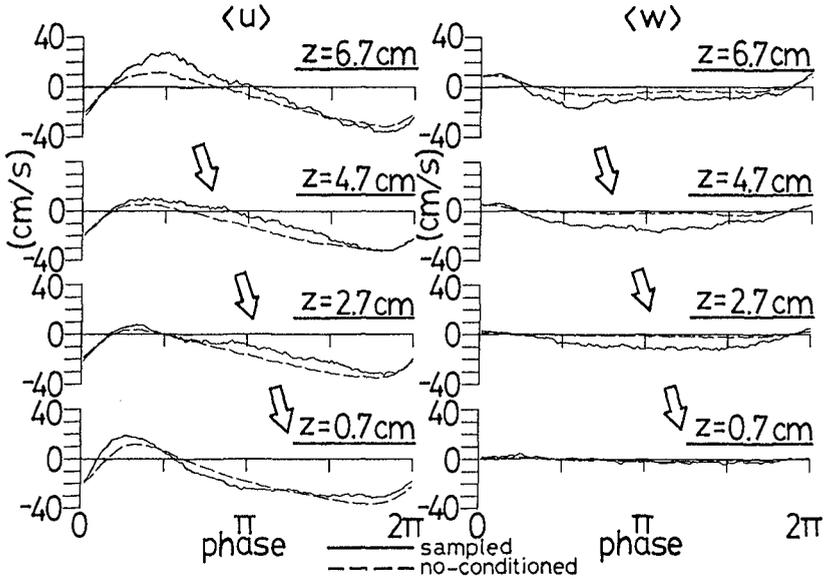


Fig. 8 Distributions of phase-averaged velocity $\langle u \rangle$ and $\langle w \rangle$ for conditionally sampled and no-conditioned records.

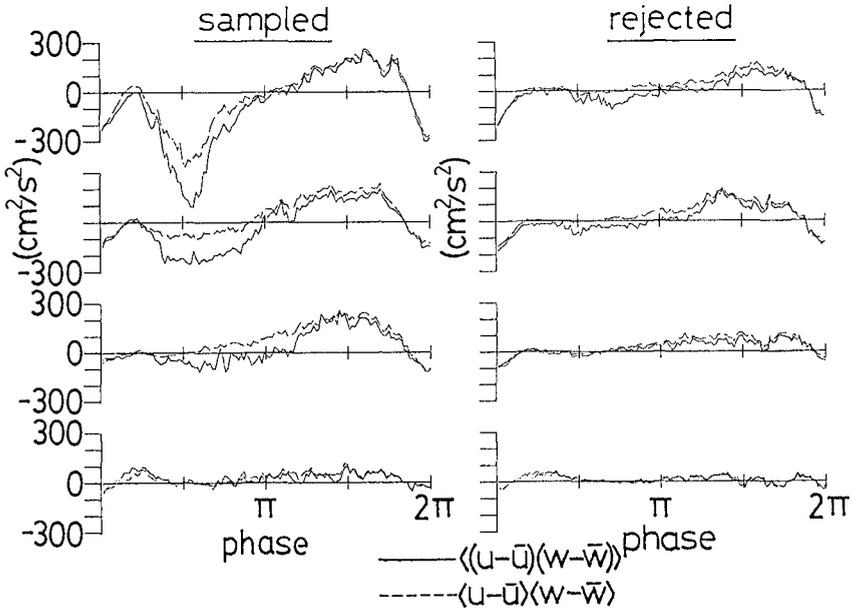


Fig. 9 Distributions of phase-averaged momentum flux for conditionally sampled and rejected records.

the left and the right figures, we can clearly see that the momentum transport of the sampled records considerably increases as a whole.

The broken lines, on the other hand, represent the product of the phase-averaged horizontal and vertical velocity fluctuation, $\langle u-\bar{u} \rangle \cdot \langle w-\bar{w} \rangle$. The difference between the solid and broken lines corresponds to the turbulent momentum transport. The fact that the difference becomes large in the highly turbulent regions means that the production of an appreciable magnitude of the Reynolds stress accompanies the evolution of the oblique eddies.

In summary, all of these results are consistent with the following conceptual model based on the obliquely descending eddies. That is, the eddies transport the intensive turbulence obliquely downward with the large onshoreward momentum at the upper layer of the water, and then brings the highly intermittent turbulent velocity and pressure fluctuation to the bottom.

Considering these results obtained through the fixed-bed experiments, we can easily expect that this eddy action to the bottom plane causes sediment suspension in a wide extent of the surf zone. To obtain more definite results on the relationship between the eddy action and resultant sediment suspension, we have made movable-bed experiments, as described as follows.

4. Relation between the Obliquely Descending Eddies and the Sediment Suspension

4.1 On-offshore distribution of the mean sediment concentration close to the bed

The on-offshore distributions of the time-averaged concentration \bar{c} both at $z=0.5\text{cm}$ located within the bedload layer and at $z=1.0\text{cm}$ just above the layer are shown in the bottom of Fig. 10. The values indicated were based on the data measured during the time interval of about 70s from the instant when the wave set-up was established until the time when the sediment ripples were generated. The initial bottom surface configuration was smoothed for each run.

The upper figure, on the other hand, shows the distribution of the bottom turbulence intensity w' and of the amplitude of orbital velocity u normalized by the linear long wave speed at the breaking point. These values of w' and u are those measured at $z=0.7\text{cm}$ in the fixed-bed experiments. Unlike for the values of the turbulence intensity indicated in Figs. 4 and 7, the averaging time for the numerical filtering of the raw data is the time of the incident wave period, because the vertical component of the near-bottom velocity contains negligibly small magnitude of the orbital velocity as compared with the turbulence. Hence w' -values so obtained include almost all contribution of the oblique eddy motion.

Comparing the shape of these distributions each other, we can recognize that the sediment concentration level \bar{c} in the bed load layer correlates well to u_m , while \bar{c} above the layer varies in accordance with the bottom turbulence inten-

sity w' . This result suggests the close relationship between the sediment suspension and the action of the oblique eddies.

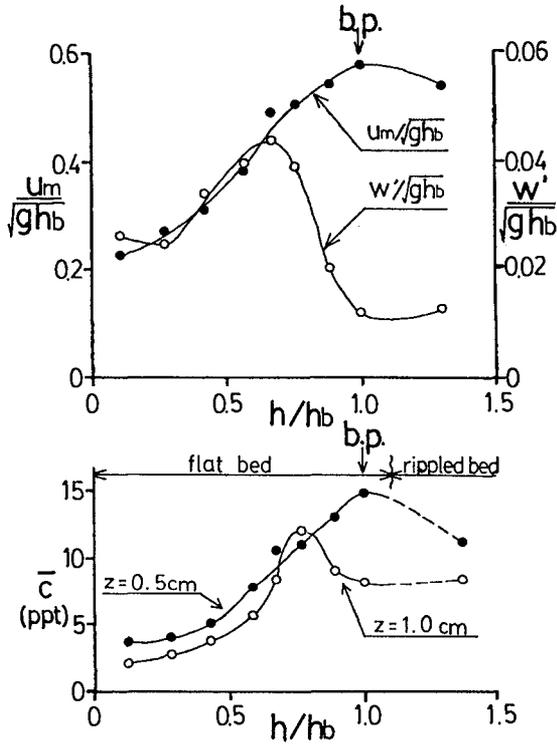


Fig. 10 Spatial distributions of mean concentration \bar{c} at $z=0.5$ and 1.0cm and those of bottom turbulence intensity w' and amplitude of orbital velocity u_m .

4.2 Visualization of the eddy evolution and the resultant occurrence of sediment suspension

To confirm this relationship from a visual aspect, we have taken some photos from a side wall of the wave channel.

Photo 1 and Figure 11 are examples showing a sequence of pictures and their tracings around the velocity measuring section taken successively by a motor-driven 35 mm camera. The vertical rod shown in the photos is the stem of the sensor to measure the sediment concentration. The instants taking the photos are indicated by the arrows denoted by the alphabets of A to H in Fig. 12, where both the water surface elevation η and the bottom sediment concentration c at $z = 1.0 \text{ cm}$ at the velocity measuring section are represented. The records shown in Fig. 12 were those obtained immediately after the starting of the wave generator; thus the first

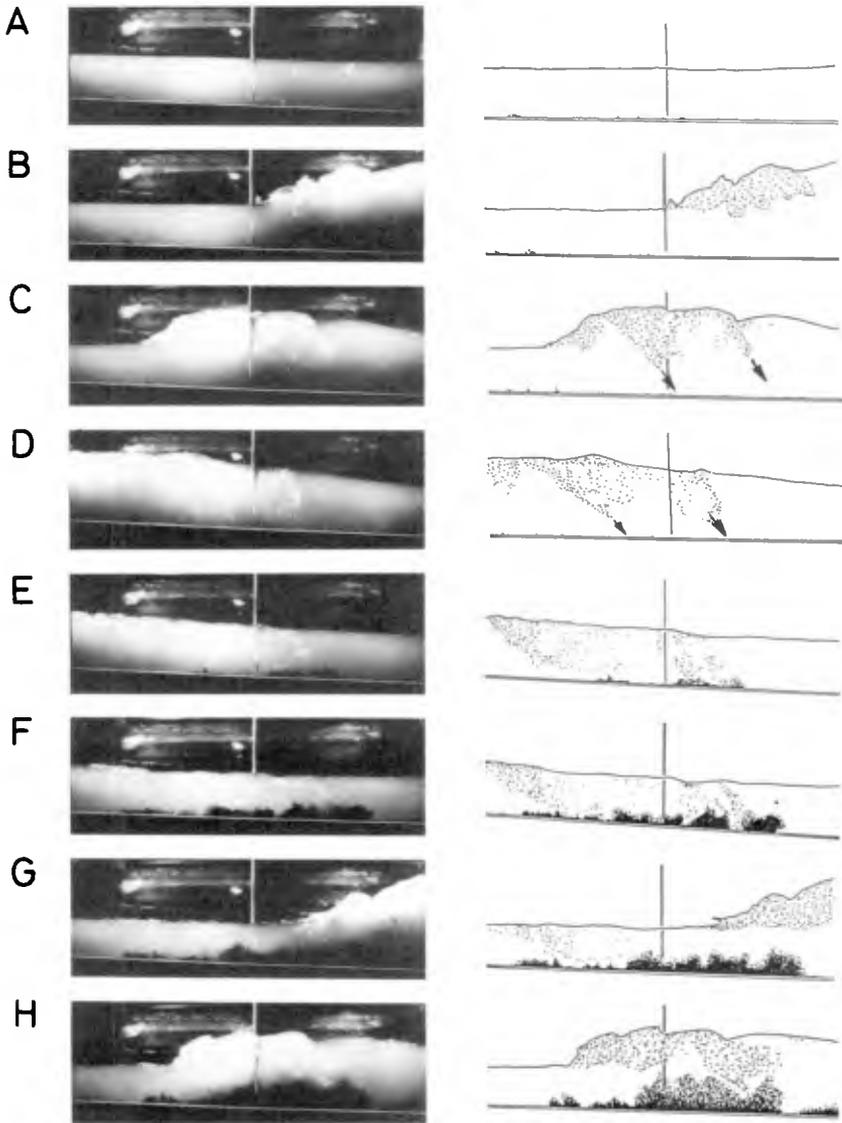


Photo 1 Development of obliquely descending eddies and sediment suspension.

Fig.11 Tracing of Photo 1.

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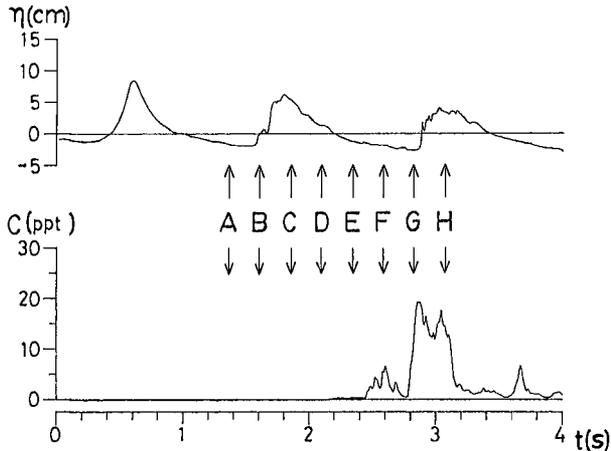


Fig.12 The instants taking the photos and the time histories of η and c .

wave indicated in the figure is a non-breaking wave. The figure indicates that the sediment concentration increases not under the first non-breaking wave, but around the phase of the trough of the second breaking wave.

The brighter regions of the photos indicate the air bubbles involved in the eddies, which are sketched in Fig.11 by the speckles. The dark regions close to the bottom in E to H, on the other hand, indicate the suspended sediment cloud. From these photos and figures, we can describe the process of the eddy development and associated sediment suspension as follows.

- (A)..... No appreciable sediment suspension can be found, though there exists the sediment movement due to sheet-flow.
- (B)..... The horizontal eddies grow to appear in front of the bore crest.
- (C)..... The eddies develop into the obliquely descending eddies behind the wave crest. The sediment suspension, however, cannot be found yet at this stage.
- (D)..... The oblique eddies stretch toward the bottom as indicated by the arrows.
- (E)..... Immediately after the instant when the eddies hit the bottom surface, the sediment suspension appears in the vicinity of the bottom.
- (F)..... A large amount of the sediment is lifted up into suspension by the eddy action, increasing the concentration value as shown in Fig. 12.
- (G)..... The sediment clouds migrate offshore with the backwash of the bore. And then they are convected upward by the heaving orbital motion preceding to the passage of the wave crest.
- (H)..... The subsequent oblique eddies evolve toward the bottom.

4.3 3-D structure of suspended sediment cloud in the surf zone

The dependence of the sand suspension on the oblique eddies mentioned above may be confirmed more quantitatively by investigating the spatial structure of the suspended sand cloud; i.e., if the oblique eddies lift up the bottom sediment into suspension, the resultant sediment cloud will show the strong three-dimensionality with the horizontal scale comparable to that of the diameter of the eddies. So to investigate the spatial structure of the sediment cloud close to the bottom, we have measured simultaneously the sediment concentration c_1 and c_2 by the two probes aligning several distances apart in the direction perpendicular to the plane of the orbital wave motion.

Figure 13 shows the time histories of the water surface elevation η as well as c_1 and c_2 both outside and inside the surf zone. The values of c_1 and c_2 in this figure are those measured by the probes manipulated to keep their heights at 1.5cm above the bottom and separated with 4cm each other in the horizontal cross-channel direction.

The values of c_1 and c_2 outside the surf zone well correlate to each other and exhibit strong dependance on the phase variation with the two marked peaks within the one

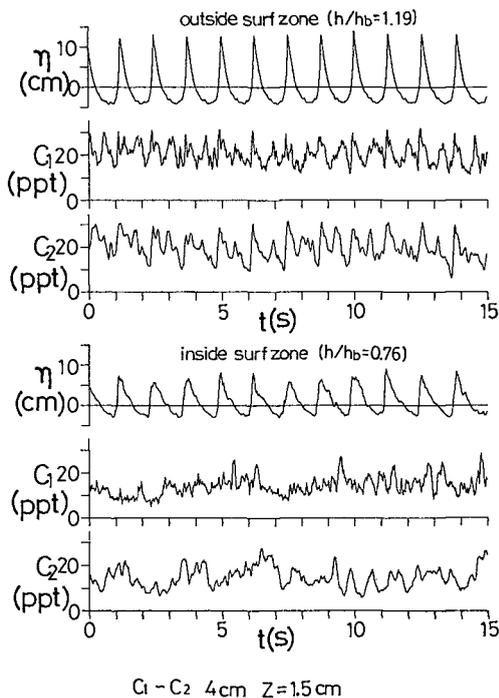


Fig. 13 An example of the time histories of η as well as c_1 and c_2 both outside and inside the surf zone.

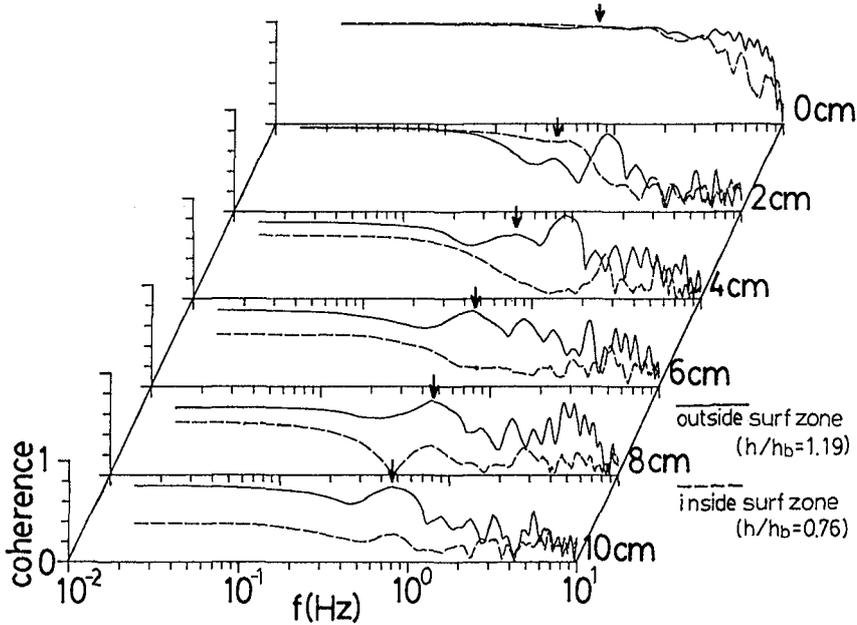


Fig.14 Coherence between bottom concentration records for several probe spacings.

wave period, though the high frequency component is slightly lost in c_2 record as compared with c_1 , because the frequency response of the probe for c_2 is lower than that for c_1 . The concentration variation inside the surf zone, on the contrary, shows no definite phase-dependance, and the correlation between these two records is rather poor. These facts suggest that the sediment suspension within the surf zone is mostly governed by the turbulence, especially by the obliquely descending eddies, generated with the wave breaking, though, in the outside of the surf zone, the wave fluid motion is a unique action to the sediment suspension, which is nearly two-dimensional phenomenon.

Then to investigate the correlation structure between c_1 and c_2 more quantitatively, we have calculated the coherence between c_1 and c_2 . The results are shown in Fig.14, where the coherence values with the horizontal axis of frequency f . The data used are those measured for the initially smoothed bed during the time interval of about 80 s from the instant when the wave set-up was established to the time when the sediment ripples were generated.

The coherence outside the surf zone denoted by full lines retain high values even for the large probe spacings, demonstrating two-dimensionality of sediment suspension outside the surf zone. On the contrary, the coherence within the surf zone represented by broken lines decreases rapidly

with the probe spacings. The decrease is most noticeable in the frequency range from 0.8 to 5 Hz which is corresponding to the period of the incident waves and its higher harmonics and is comparable to the frequency of the eddy occurrence. This means that the sediment suspension inside the surf zone occurs in a quite spot-like manner and is consistent again with the proposed model of sediment suspension by the oblique eddies.

5. Conclusions

The results and conclusions obtained in the present study can be summarized as follows.

1) "Obliquely descending eddies" bring highly intermittent turbulence to the bottom with the large onshoreward momentum and thus essentially characterize the turbulent flow field in the surf zone.

2) The eddies hit the bottom and then lift up the sediment into suspension, yielding the spot-like sediment cloud in accordance with the 3-dimensional eddy structure.

Acknowledgment

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